Analysis of L-Probe Fed-Patch Microstrip Antennas in a Multilayered Spherical Media

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Abstract—In this paper, the L-probe fed-patch microstrip antennas in a multilayered spherical media are analyzed based on the method of moments (MoM). Firstly, the structures of microstrip antennas in three-layered spherical media fed by L-probe are established as similar as the plane case. Then, the method of moments is employed to solve the mixed potential integral equation (MPIE). Numerical results for the input impedance, radiation patterns are showed good agreement with ones come from commercial software FEKO. Finally, the effects of all kind of geometric parameters on the characteristics of antenna are investigated for the next design purpose.

I. INTRODUCTION

As we know, the microstrip antenna is a kind of resonance antenna and the narrow bandwidth is the major weakness. Much effort has been done to widen the bandwidth of microstrip in plane case in past decades, including impedance matching technology [1] and using thick substrates [2]. The feeding approach employing L-probe is an effective method to ameliorate bandwidth, which is especially convenient for the coplanar case. The microstrip antenna with an L-probe in planar case has been studied in [3]-[5]. A modified L-probe fed microstrip antenna has been presented in [6].

Compared with in planar case, only a few literatures relate to the microstrip antenna with L-probe in coplanar case [7]. The conformal microstrip antennas have a major advantage that they have a wider angular radiation pattern than the planer counterparts. It is worthwhile to research the characteristics of the antennas with L-probe in conformal case.

In this paper, the L-probe patch antenna in a multilayered spherical media is studied by employing the method of moments (MoM). The structure of antenna is obtained with the aid of commercial software FEKO, which is divided into a mesh of triangular patches. The surface current density can be calculated by solving the mixed potential integral equation (MPIE) based on the Rao-Wilton-Glisson (RWG) basis function. Characteristics including input impedance, radiation pattern are presented, which agree well with those from FEKO. Finally, extensive numerical results of antenna parametric study are showed in order to provide some design information.

II. THEORY

A. Structure of Antenna

The structure of antenna has three layers as shown in Fig.1. The innermost layer is a perfectly conducting spherical core, which can be modeled assuming the permeability tends to zero and permittivity tends to infinity [8]. The outermost layer assumed to be free space. The rectangular patch locates at interface of medium layer and free space and the conformal L-probe positions at the medium layer. Detailed parameters are shown in Fig.1. In this paper, nonmagnetic materials have been considered, that is $\mu_r = 1$.

![Figure 1. Structure of antenna with L-probe](image)

B. Mixed potential integral equation

In this paper, MoM is employed to solve the problem because of its accuracy and efficient. In order to make analysis process simplification, the L-probe can be treated as an L-trip [9], which is connected to the PEC spherical core. A delta gap voltage source has been placed at the base of the feeding trip. The tangential electric field can be achieved employing the mixed potential integral equation:
\[ E'_i = -j\omega\mu_i \int\int G_i \cdot J_i(r')ds' - \nabla \int\int G_e \cdot J_i(r')ds', \quad (1) \]

where \( \mu_i \) represents the permeability of the layer of field point. \( s' \) is the surface area of the antenna conducting body except the ground sphere as shown in Fig.2, and \( J_i(r') \) is the current density of the surface area. The contribution of the ground sphere has been incorporated in the Green’s functions including magnetic vector potential Green’s function \( \vec{G}_m \) and electric scalar potential Green’s function \( G_e \).

Solving the dyadic Green function in spherically multi-layered media is indispensable in order to make the analysis results accurate. A general expression of the dyadic Green’s function for the electromagnetic fields in an arbitrary multi-layered medium has been achieved in [10]. Further development of Green’s function has been proposed in [8], where the electric field Green’s function was transformed to the mixed potential Green’s function. Take \( G_e \) component for example:

\[
G_e = \frac{j\omega\mu_i}{4\pi} \sum_{n=0}^{n+1} 2n+1 \times \\
\left[ \frac{\partial h_e^{(2)}(k,r)}{\partial k} \left( A_{mn}^{(2)} \delta h_e^{(2)}(k,r) + B_{mn}^{(2)} \delta j_{e}(k,r') \right) + \delta j_{e}(k,r) \left( C_{mn}^{(2)} \delta h_e^{(2)}(k,r) + D_{mn}^{(2)} \delta j_{e}(k,r') \right) \right] \cdot P_\gamma(\cos \gamma) \quad (2)
\]

meaning of the denotations can be found in [8]. From the expressions, we can see that the above expression is an infinite summation of spherical harmonics. In order to accelerate the convergence of the infinite series, asymptotic extraction approach has been adopted in [11]-[12].

The surface area of the antenna conducting body has been divided into a mesh of triangular patches as shown in Fig.2, and the RWG triangular function is the most appropriate basis function [13]. Then, we expand the current \( J_i(r') \) via a sum of weighted basis functions:

\[ J_i(r') = \sum_{n=1}^{N} a_n f_n(r') \quad (3) \]

and the method of Galerkin has been adopted in this paper. Every element of the impedance matrix can be obtained by the expression below:

\[
Z_{mn} = \int\int \int \left[ j\omega\mu_i \cdot f_m(r) \cdot \vec{G}_i \cdot f_n(r') + (\nabla \cdot f_m(r))G_e(\nabla \cdot f_n(r')) \right] ds'ds' \quad (4)
\]

where \( m \) and \( n \) refer to the test and expansion function.

The electric field can be expressed as:

\[
E_e = \int\int j\omega\mu_i \vec{G}_e \cdot f_i(r') + \nabla(\nabla \cdot G_e \cdot f_i(r')) ds' \quad (5)
\]

A problem that needs special attention is the connection of the feeding trip to the perfectly conducting sphere. The delta gap voltage source has been placed at the connection, and more detailed analysis can be found in [14].

C. Numerical results

An L-probe fed-patch microstrip antenna in multilayered...
From above figures, we can see that the calculated result of input impedance show good well agreement with the one from FEKO, and a small error is observed between far-field radiation pattern calculated and FEKO. The main reason is that the meshes of the two antennas are different and only surface area of antenna is meshed in this paper, while the whole antenna is meshed in FEKO. The error can be ignored that is acceptable in engineering designs. Fig. 6 shows the SWR of the antenna calculated by this paper and FEKO. From the figure, we can see the antenna attains an about 1.7 GHz impedance bandwidth (from 3.95 GHz to 5.65 GHz, $SWR \leq 2$), which is similar to that in planar case [2], and wider than that of antenna fed by straight probe.

**III. PARAMETRIC STUDY**

In this section, we will investigate the variation of the antenna characteristics with the antenna parameters including $W_1$, $H_1$, $L_1$, $L_2$ and $D$. Because of the limited paper space, only the input impedances are shown, and the far-field radiation pattern and SWR of the antennas are not shown in this paper.

The variation of the input impedance with the length of patch $W_1$ is illustrated in Fig. 7. From the figure, we can see that input resistance decreases with increasing $W_1$ and the input reactance decreases before the resonant frequency and increases after the resonant frequency with increasing $W_1$.

The variation of the input impedance with the patch width $H_1$ is shown in Fig. 8. It can be observed that $H_1$ determined the resonant frequency of the antenna. The resonant frequency becomes high with decreasing $H_1$. Fig. 9 and Fig. 10 show the variation of the input impedance with the vertical length $L_1$ and the horizontal length $L_2$ of the L-strip, respectively. We can see from the figures that both the input resistance and reactance increase with $L_1$, $L_2$. At the resonant frequency, the input resistance increases clearly with increasing $L_1$, and the reactance increases clearly with increasing $L_2$.

The effect of changing the position $D$ of the L-strip from the patch is studied as shown in Fig. 11. $D > 0$ denotes that the L-strip is partly covered by the patch, and $D < 0$ denotes the strip is covered by the patch in whole. We can see that both the resistance and reactance change slightly when $D \geq 0$, and decrease clearly when $D < 0$. 
Figure 10. Input impedance versus frequency for different $L_2$
$W = 30\text{mm}$, $H = 25\text{mm}$, $L_1 = 4.95\text{mm}$, $h = 6.6\text{mm}$ and $D = 2\text{mm}$

Figure 9. Input impedance versus frequency for different $L_1$
$W = 30\text{mm}$, $H_1 = 25\text{mm}$, $L_2 = 10\text{mm}$, $h = 6.6\text{mm}$ and $D = 2\text{mm}$

Figure 11. Input impedance versus frequency for different $D$
$W = 30\text{mm}$, $H = 25\text{mm}$, $L_1 = 4.95\text{mm}$, $L_2 = 10\text{mm}$ and $h = 6.6\text{mm}$

IV. Conclusion

In this paper, the L-probe patch antenna in a multilayered spherical media is studied by employing the method of moments. The RWG function is regarded as bases function and test function for solving the mixed potential integral equation. Numerical results for input impedance, far-field radiation pattern are presented and verified by comparison with the results from commercial software FEKO. The effects of various parameters on the characteristics of the L-strip patch microstrip antenna have been investigated. From results obtained, we can see that the input impedance is sensitive to $W$, $L_1$ and $L_2$. The resonant frequency is determined by the parameter $H_1$. The effects of various parameters here are similar to that in planar case.

REFERENCES